

# Control of Optical Performance in Infrared Region for Vanadium Dioxide Films Layered by Amorphous Silicon

H. Kakiuchida · P. Jin · M. Tazawa

Published online: 5 March 2009  
© Springer Science+Business Media, LLC 2009

**Abstract** The multilayer structure of amorphous silicon (*a*-Si) with vanadium dioxide (VO<sub>2</sub>) that possesses excellent thermochromism in the infrared range was formed by reactive magnetron sputtering. Then, the physical properties such as thermochromic performance, crystalline structure, and film morphology were investigated. As a result, the *a*-Si layer was found to flexibly control the optical-spectral variation due to thermochromism produced in the VO<sub>2</sub> layer over a wide infrared range. The multilayer structure of VO<sub>2</sub> between *a*-Si layers is an effective way to design various infrared thermochromic devices, since the refractive index of *a*-Si can be widely varied by oxidation.

**Keywords** Amorphous silicon · Infrared radiation control · Magnetron sputtering · Semiconductive-to-metallic phase transition · Thermochromism · Vanadium dioxide

## 1 Introduction

Vanadium dioxide (VO<sub>2</sub>), which exhibits a semiconductor-to-metal phase transition, has excellent thermochromic performance, that is, a large variation in optical constants particularly at near-infrared and longer wavelengths [1,2]. Since the transition temperature is at approximately 68 °C, and additionally it can be reduced to the vicinity of room temperature by a trace of dopant [3,4], such thermochromic function is applicable to solar-radiation control for seasonably comfortable houses [5]. In summer, the intensity of solar radiation into a room through a window decreases, while it increases

---

H. Kakiuchida (✉) · P. Jin · M. Tazawa  
Materials Research Institute for Sustainable Development, National Institute  
for Advanced Industrial Science and Technology (AIST), 2266-98 Anagahora,  
Shimoshidami, Moriyama, Nagoya 463-8560, Japan  
e-mail: h.kakiuchida@aist.go.jp

in winter. Since the transmittance of  $\text{VO}_2$  is controlled automatically by a variation in ambient temperature due to seasonal changes, no switching devices are needed.

Optical spectra for the low- and high-temperature phases of  $\text{VO}_2$ , which correspond to semiconductive and metallic phases, respectively, are still desired to be more highly controllable at infrared wavelengths. We have earlier proposed several multilayer structures composed of  $\text{VO}_2$  and transparent materials such as  $\text{TiO}_2$ ,  $\text{ZrO}_2$ , etc. [5, 6]. However, the refractive index of  $\text{VO}_2$  is so large that optical materials with a higher refractive index, preferably over three, are desired when they are layered with  $\text{VO}_2$  for better controlling the spectral change due to thermochromism.

This study focuses on the multilayer structure of amorphous silicon ( $a\text{-Si}$ ) with  $\text{VO}_2$  film, because  $a\text{-Si}$  film has several advantages: (i) refractive index higher than that of  $\text{VO}_2$ , (ii) high controllability of refractive index by oxidation [7], and (iii) manageability in deposition technique and cost. From the viewpoint of practical fabrication of this multilayer structure, material suitability of  $a\text{-Si}$  film for  $\text{VO}_2$  film is indispensable. The under- (over-)coatings of the  $\text{VO}_2$  layer by amorphous materials may affect the physical properties of  $\text{VO}_2$  such as the stability of crystalline growth, film morphology, and thermochromic functionality, because the crystalline structure of the vanadium-oxide system ( $\text{VO}_n$ ;  $0 < n < 2.5$ ) [8] is sensitive to under- (over-)coated materials. For example, Muraoka et al. [9] have examined the electrical-resistance variation due to the phase transition of  $\text{VO}_2$  film deposited on a  $\text{TiO}_2$  substrate with various crystal faces, and then they have reported the correlation between the phase-transition temperature and lattice parameter of  $\text{VO}_2$ . On the other hand, Kato et al. [10] have shown that the phase-transition behavior is strongly affected by the thickness of a zinc oxide ( $\text{ZnO}$ ) layer as an undercoat of  $\text{VO}_2$  film. According to the previous results, the phase transition or thermochromism strongly depends on the lattice (mis)match with crystalline materials under the  $\text{VO}_2$  film. Thus, the potential of amorphous materials for stable deposition of  $\text{VO}_2$  film is considerable, since the amorphous structure has no specific crystal lattice.

In this study, the feasibility to layer  $\text{VO}_2$  film with  $a\text{-Si}$  film was studied with regard to material suitability and optical design. By using radio frequency (rf)-reactive magnetron sputtering, which can possibly disturb the stability of  $\text{VO}_2$  crystalline growth but must be used for practical reasons, the multilayer structure composed of  $\text{VO}_2$  film sandwiched by a pair of  $a\text{-Si}$  layers (SVS structure) was formed, and then the crystalline structure, film morphology, and optical properties of  $\text{VO}_2$  and  $a\text{-Si}$  were investigated.

## 2 Experimental and Estimation Procedures

The films were deposited by rf-reactive magnetron sputtering using metal targets.  $\text{VO}_2$  film was deposited by sputtering a vanadium (V) target in argon (Ar) and oxygen ( $\text{O}_2$ ) gases under flow rates of 32 sccm and 1.8 sccm, respectively, with a total pressure of 0.6 Pa, rf power of 100 W, and substrate temperature of 600 °C. On the other hand,  $a\text{-Si}$  film was deposited by sputtering a silicon (Si) target in Ar gas at a flow rate of 32 sccm, 0.6 Pa in total pressure, 120 W in rf power, and 600 °C in substrate temperature. Before sputtering was initiated, the vacuum chamber was evacuated to  $2 \times 10^{-6}$  Pa. The  $\text{VO}_2$

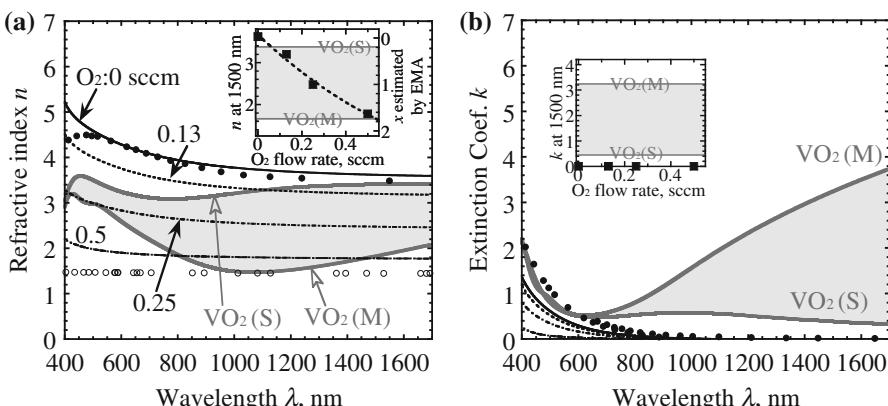
films sandwiched by a pair of  $a\text{-Si}$  layers (SVS structure) at several thicknesses, which have the effect of controlling the optical spectrum, were designed as optical multilayers by coordination of the sputtering procedure.

To determine the optical constants and deposition rates of the present  $\text{VO}_2$  and  $a\text{-SiO}_x$  films, single-layered samples had been preliminarily prepared and the thickness and optical constants were analyzed by conventional ellipsometry. The measurements were carried out using an ellipsometer (J. A. Woolam Co., Inc., M2000).

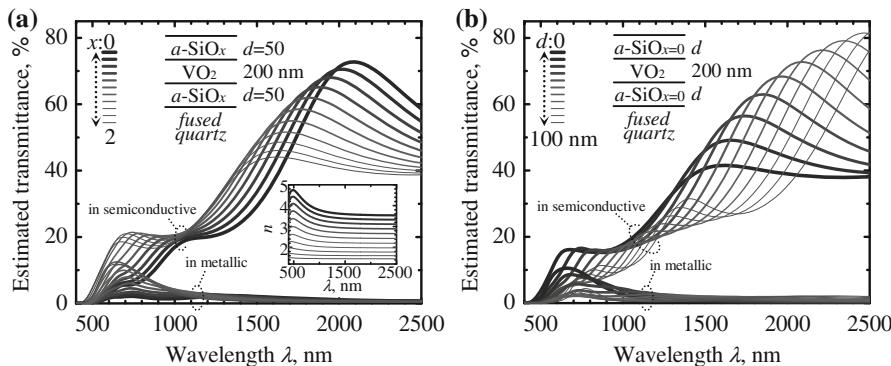
The thermochromic behavior of the  $\text{VO}_2$  film between  $a\text{-Si}$  layers was examined at visible and near-infrared wavelengths by a spectrophotometer (Hitachi High-Technologies Corp., U4100), and at infrared wavelengths by an FTIR (Perkin Elmer Inc., GX1P). The morphology of the SVS structure was observed by a scanning electron microscope (SEM) (Hitachi High-Technologies Corp., S-4300) and an atomic force microscope (AFM) (Veeco Instruments Inc., MultiMode). The crystalline information was obtained by X-ray diffractometry (XRD) in the  $\theta\text{--}2\theta$  mode with  $\text{CuK}\alpha$  (Rigaku Corp., RINT2000).

### 3 Estimation, Results, and Discussion

Figure 1 shows the optical constants of  $a\text{-SiO}_x$  films prepared by reactive sputtering of a silicon target in various  $\text{O}_2$  atmospheres, where the other sputtering conditions were the same as in  $a\text{-Si}$  deposition. The reported data for  $a\text{-Si}$  and  $a\text{-SiO}_2$  are also plotted in the figure [11, 12]. The insets show near-infrared optical constants plotted against the flow rate of  $\text{O}_2$  during the sputtering process. The right-hand ordinate of the inset in Fig. 1a indicates the oxidation number,  $x$ , which was estimated through the



**Fig. 1** Optical constants of  $a\text{-SiO}_x$  (black curves), deposited by sputtering a Si target at various  $\text{O}_2$  flow rates, where closed and open circles represent the reported values of  $a\text{-Si}$  and  $a\text{-SiO}_2$ , respectively [11, 12]. Insets show the optical constants and oxidation number of  $a\text{-SiO}_x$  at a wavelength of 1500 nm as a function of  $\text{O}_2$  flow rate, where  $x$  was estimated by EMA from the optical constants of  $a\text{-Si}$  and  $a\text{-SiO}_2$ , as detailed in the text. The values of  $\text{VO}_2$  change reversibly with temperature by a semiconductive-to-metallic (S-to-M) phase transition, as expressed by gray curves and areas, where the values of S and M phases were measured at 30 °C and 90 °C, respectively

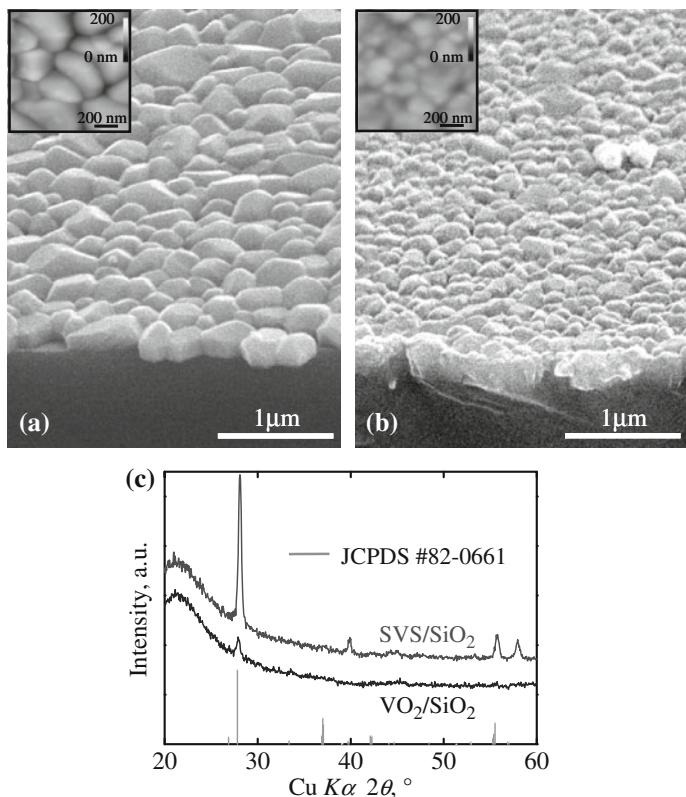


**Fig. 2** Optical transmittances of SVS structures on fused quartz at various (a) oxidation numbers,  $x$ , and (b) thicknesses of  $a\text{-SiO}_x$  layers,  $d$ , which were estimated from optical constants of semiconductive and metallic  $\text{VO}_2$  and  $a\text{-SiO}_x$ , as shown in Fig. 1. The transmittances were simulated by using refractive indices of  $a\text{-SiO}_x$  at several  $x$  as shown in the inset, and those of  $\text{VO}_2$  (S and M) as shown in Fig. 1. The layer structures are schematically shown in the figures

Bruggeman effective medium approximation (EMA) [13] from the reported optical constants of  $a\text{-Si}$  and  $a\text{-SiO}_2$ . The refractive index of  $a\text{-SiO}_x$  decreases monotonically with an increase in the flow rate of  $\text{O}_2$ , and it covers the values of both semiconductive and metallic  $\text{VO}_2$ . On the other hand, the extinction coefficient of  $a\text{-SiO}_x$  is approximately zero in the near-infrared range, although a shoulder of the optical absorption edge appears. This suggests high controllability of thermochromic behavior by introduction of  $a\text{-SiO}_x$  layers multilayered with  $\text{VO}_2$ . For example, the transmittance spectra of an SVS structure including a 200 nm thick  $\text{VO}_2$  layer were estimated from the optical constants of  $\text{VO}_2$  and  $a\text{-SiO}_x$ , as shown in Fig. 2a, b. The refractive indices of  $a\text{-SiO}_x$  used for the transmittance estimations are shown as the inset in Fig. 2a, which were obtained by the aforementioned EMA. Particularly in the near-infrared range, the transmittance change due to thermochromism produced in the  $\text{VO}_2$  layer is well controlled by the thickness and/or oxidation number of  $a\text{-SiO}_x$  layers.

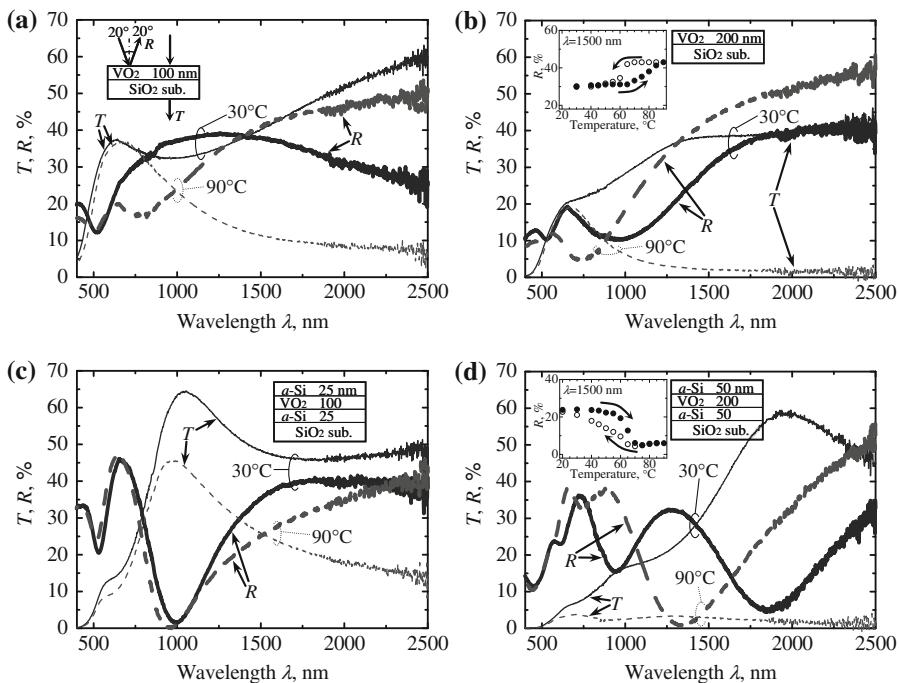
Figure 3a, b show SEM images of cross sections of the  $\text{VO}_2$  single- and SVS multilayers on a fused-quartz substrate, respectively. The insets show close-up views of their morphologies as measured by AFM. Crystalline-like grains with submicron size appear in both samples. The grains observed in the SVS multi-layer sample (Fig. 3b) are, on the average, smaller and more closely packed than those in the  $\text{VO}_2$  single-layer sample (Fig. 3a). Figure 3c shows XRD patterns of the  $\text{VO}_2$  single- and SVS multilayers on a fused-quartz substrate in the range between  $20^\circ$  and  $60^\circ$  in  $2\theta$ , where solid bars at the bottom of the graphs are the peaks identified by crystalline  $\text{VO}_2$  [14]. The peak at around  $28^\circ$  in  $2\theta$ , which corresponds to (011) interplanar spacing of monoclinic  $\text{VO}_2$ , appears in both samples. The peak intensity seems to be larger in the SVS structure than in the  $\text{VO}_2$  single layer. The information for the film morphology and crystalline quality, as aforementioned, suggests that  $\text{VO}_2$  can be formed even in the SVS structure and an  $a\text{-Si}$  layer does not decrease the stability of the  $\text{VO}_2$  formation.

Since  $a\text{-Si}$  has a large refractive index and a small extinction coefficient at near-infrared wavelengths, the SVS structure is expected to provide excellent controllability



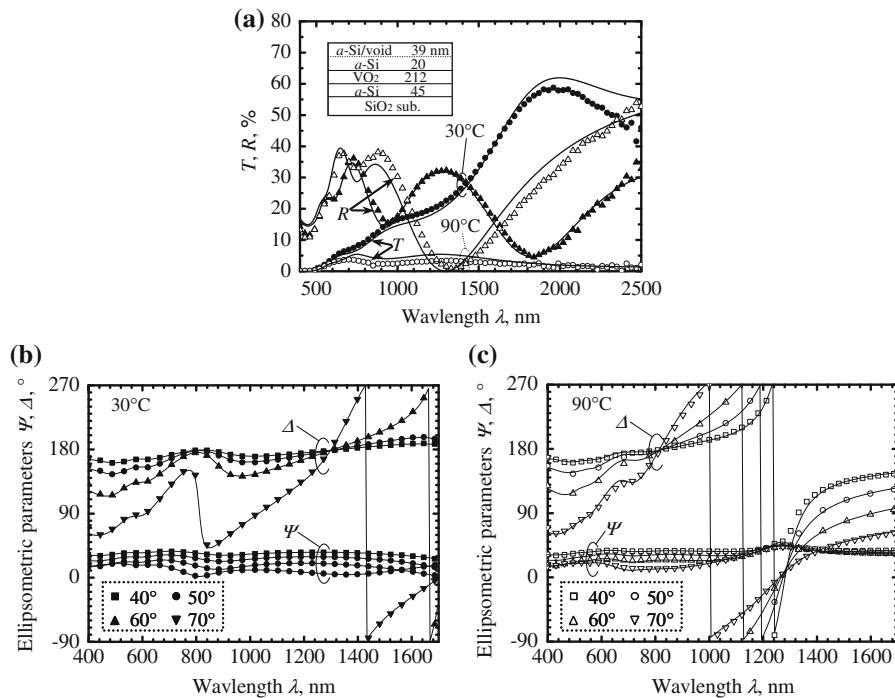
**Fig. 3** SEM images of cross sections of (a) VO<sub>2</sub> single- and (b) SVS multilayer films on fused quartz. Insets show surface configurations measured by AFM, and c XRD patterns, where the peak around 28° is attributed to (011) monoclinic VO<sub>2</sub> [14]

of the change in the optical spectrum due to thermochromism of VO<sub>2</sub>. The transmittance ( $T$ ) and reflectance ( $R$ ) of VO<sub>2</sub> single- and SVS multi-layers with different thicknesses on the fused-quartz substrate are plotted against wavelength in Fig. 4a–d.  $R$  at a wavelength of 1500 nm is plotted against temperature in the insets in Fig. 4b, d. The layer structures are schematically shown in the respective figures. For  $T$  and  $R$  measurements, the angles of incidence are 0° and 20°, respectively, as schematically shown in Fig. 4a, and the incident light was unpolarized and the detected light was not selective with regard to polarization. As a whole, thermochromic performance clearly appears in both VO<sub>2</sub> single- and SVS multi-layer samples, which bears direct evidence from the viewpoint of optical performance: the  $a$ -Si layer does not result in a decline of the crystalline quality of VO<sub>2</sub> and consequent thermochromism. The changes in  $T$  and  $R$  due to thermochromism seem to be marked rather in the near-infrared range. While the optical spectra are simple for the VO<sub>2</sub> single-layer structure, as shown in Fig. 4a, b, the optical spectra are drastically transfigured by the sandwich of VO<sub>2</sub> by  $a$ -Si layers, as shown in Fig. 4c, d. The changes in transmittance and reflectance due to thermochromism around specific wavelengths can be flexibly expanded or reduced.

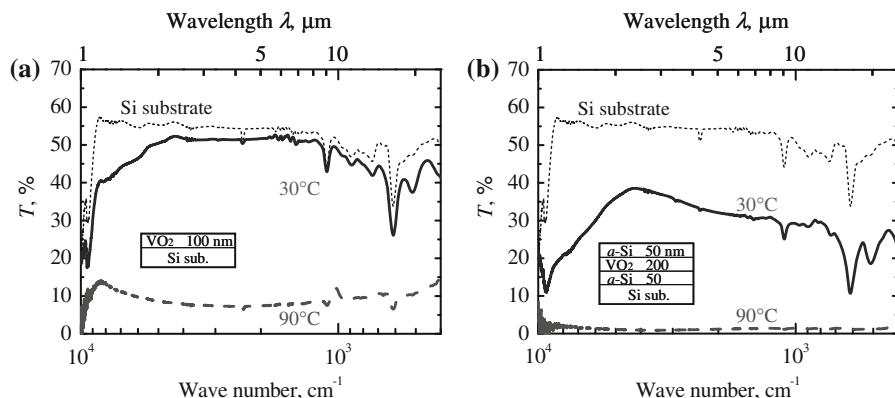


**Fig. 4** Transmittance ( $T$ ) and reflectance ( $R$ ) spectra of various film structures including thermochromic  $\text{VO}_2$  layer, where the layer structures are schematically shown in the respective figures.  $T$  and  $R$  were measured for semiconductive ( $30^\circ\text{C}$ ) and metallic ( $90^\circ\text{C}$ ) phases of  $\text{VO}_2$ .  $R$  at a wavelength of  $1.5\ \mu\text{m}$  is plotted against temperature in (b) and (d). Insets show the reflectance at a wavelength of  $1500\ \text{nm}$  against temperature, where the closed and open circles are the results measured in rise and drop in temperature, respectively. The kinks in the spectra at  $850\ \text{nm}$  result from the change of detectors between visible and infrared ranges

However, these measured spectra deviate from the expected ones as shown in Fig. 2. To discuss the origin of these deviations, the SVS multilayer structure whose spectra are shown in Fig. 4d was ellipsometrically analyzed. Figure 5a shows transmittance ( $T$ ) and reflectance ( $R$ ) spectra together with the simulations from the ellipsometric analyses. The analyzed multilayer structure is shown in the same figure. In the present analyses, the EMA layer composed of  $a\text{-Si}/\text{void}$  mixture was employed for expressing the rough surface of the top  $a\text{-Si}$  layer [15], which was modeled on the basis of the observed cross section of the film as shown in Fig. 3b. The measured ellipsometric parameters ( $\Psi$ ,  $\Delta$ ) at both  $30^\circ\text{C}$  and  $90^\circ\text{C}$  agree with the fits, as shown in Fig. 5b, c, respectively, and the measured  $T$  and  $R$  were consistently reproduced by the present analyses, as shown in Fig. 5a. Accordingly, the top  $a\text{-Si}$  layer possesses an uneven surface of  $39\ \text{nm}$  in thickness, which is different from the designed layer, while the  $\text{VO}_2$  and bottom  $a\text{-Si}$  layers are ideally formed. Therefore, such large surface roughness of the top  $a\text{-Si}$  layer is regarded as the origin of deviation of the measured transmittances from the designed ones.



**Fig. 5** Transmittance ( $T$ ), reflectance ( $R$ ), and ellipsometric parameters ( $\Psi$ ,  $\Delta$ ) of SVS multi-layer structure at 30 °C and 90 °C. (a)  $T$  and  $R$  reproduced from Fig. 4d, as indicated by closed (30 °C) and open (90 °C) symbols. Inset shows film structure analyzed by ellipsometry. Solid curves are the simulated spectra from the ellipsometric analysis, and (b) and (c) ellipsometric parameters at 30 °C and 90 °C, respectively. Symbols and solid curves represent the measurements and fits, respectively.  $\Psi$  and  $\Delta$  were measured at angles of incidence of 40°, 50°, 60°, and 70°



**Fig. 6** Infrared transmittance spectra of (a)  $\text{VO}_2$  single- and (b) SVS multilayers on crystalline silicon substrate, where the layer structures are schematically shown in the figures. Note that the spectra contain several absorptions due to vibrational modes of extrinsic materials such as  $\text{H}_2\text{O}$ ,  $\text{CO}_2$ , organic compounds, etc. Transmittances were measured for semiconductive (30 °C) and metallic (90 °C) phases of  $\text{VO}_2$

Transmittances of VO<sub>2</sub> single- and SVS multilayer samples at longer wavelengths in the infrared range are shown in Fig. 6a, b. For transmittance measurements in the far-infrared range, crystalline silicon was chosen as a substrate instead of fused quartz. Although the optical design of the present films was not specifically optimized for the far-infrared range, the results exhibit large potential for controlling thermochromic performance in the far-infrared spectrum as well as in the near-infrared region by formation of the appropriate SVS structure.

#### 4 Conclusions

This study demonstrated that layering *a*-Si with VO<sub>2</sub> is an efficient and reliable technique for controlling the change in optical spectrum due to thermochromism at near- and far-infrared wavelengths. In particular, the sandwich structure of VO<sub>2</sub> between *a*-Si layers is an effective way for extending the application range of thermochromic devices, since the refractive index and thickness of the *a*-Si film are highly controllable.

#### References

1. F.J. Morin, Phys. Rev. Lett. **3**, 34 (1959)
2. J.B. Goodenough, J. Solid State Chem. **3**, 490 (1971)
3. I. Takahashi, M. Hibino, T. Kudo, Jpn. J. Appl. Phys. **40**, 1391 (2001)
4. C.G. Granqvist, Sol. Energy Mater. Sol. Cells **91**, 1529 (2007)
5. P. Jin, G. Xu, M. Tazawa, K. Yoshimura, Appl. Phys. A **77**, 455 (2003)
6. G. Xu, P. Jin, M. Tazawa, K. Yoshimura, Sol. Energy Mater. Sol. Cells **83**, 29 (2004)
7. E. Dehan, P. Temple-Boyer, R. Henda, J.J. Pedroviejo, E. Scheid, Thin Solid Films **266**, 14 (1995)
8. C.H. Griffiths, H.K. Eastwood, J. Appl. Phys. **45**, 2201 (1974)
9. Y. Muraoka, Z. Hiroi, Appl. Phys. Lett. **80**, 583 (2002)
10. K. Kato, P. K. Song, H. Odaka, Y. Shigesato, Jpn. J. Appl. Phys. **42**, 6523 (2003)
11. H. Piller, in *Handbook of Optical Constants of Solids*, ed. by E. D. Palik (Academic Press, San Diego, 1998), pp. 571–586
12. I.H. Mallitson, J. Opt. Soc. Am. **55**, 1205 (1965)
13. D.E. Aspnes, Thin Solid Films **89**, 249 (1982)
14. JCPDS #82–0661
15. H. Fujiwara, J. Koh, P.I. Rovira, R.W. Collins, Phys. Rev. B **61**, 10832 (2000)